

The underwater light field

by Dr Derek Pilgrim

The light field describes the *flux* of propagating light. In the definitions following, angles will be expressed in radians. Directions will be specified by the zenith angle, θ , ($\theta = 0 =$ zenith; $\theta = \pi =$ nadir), and the azimuth, φ , ($0 \geq \varphi \geq 2\pi$). The *units* of the quantities involved are given in square parentheses, [].

Radiant energy, $Q(\lambda)$ [J or q = quanta]

This is the radiant energy that moves in the radiation field, indeed, the whole physical purpose of radiant propagation is to move energy from one place to another.

Radiant flux, $F(\lambda)$ [W, qs^{-1} , lumen]

Since radiation must be propagating at velocity, c , then flux of radiant energy (*i.e.* radiant power), is the fundamental quantity used to describe a radiation field. It is the time rate (d/dt) flow of radiant energy $Q(\lambda)$, and so:

$$F(\lambda) = \frac{dQ(\lambda)}{dt} \quad (\text{eqn. 1})$$

The photopic (visual) unit of radiant or luminous flux is the lumen, a measure of the power of light perceived by the human eye. It differs from radiant flux, F , in that luminous flux takes into account the varying sensitivity of the human eye to different wavelengths of light.

Radiant intensity, $I(\lambda)$ [Wsr⁻¹, $qs^{-1}sr^{-1}$, lumen sr⁻¹]

Radiant intensity is the measure of the radiant flux per unit solid angle, ω , (per

steradian, sr^{-1}). Consider a radiance source radiating a radiant flux of $F(\lambda)$ watts as in *Fig.1*. If this flux is concentrated into a cone of solid angle 1 steradian, then it has an intensity of $I \text{ W sr}^{-1}$. We may therefore write:

$$I = \frac{dF(\lambda)}{d\omega} = \frac{d^2Q}{d\omega \cdot dt} \quad (\text{eqn.2})$$

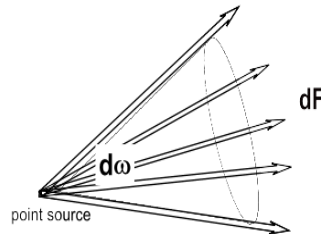


Fig.1 Radiant intensity, $I(\lambda)$

Radiance, $L(\lambda)$ [$\text{Wsr}^{-1}\text{m}^{-2}$, $\text{qs}^{-1}\text{sr}^{-1}\text{m}^{-2}$, lux]

Suppose that we now replace the radiant source in *Fig.1* with a radiance detector of infinitesimal size. If the field of view of the detector is restricted by a conical shield to an angle of 1 steradian then it will measure radiant flux in units of watts per steradian [Wsr^{-1}]. However, if the detector has finite area, dA , and is placed at the end of a tube (Gershun tube) of length l as in *Fig.2*, then the field of view of the detector will become restricted to a solid angle of $d\Omega = dA/l^2$, so that the radiant flux (power) arriving at the detector will be: $dF(\lambda)/d\Omega$.

Fig.2 Radiance, $L(\lambda)$

Since the detector has finite area, dA , then the intensity that it measures will be the radiance, $L(\lambda)$, *i.e.*:

$$L(\lambda) = \frac{d^2F(\lambda)}{d\Omega \cdot dA} \quad \text{W sr}^{-1}\text{m}^{-2} \quad (\text{eqn.3})$$

Strictly, in defining radiance, $d\Omega \rightarrow 0$, *i.e.* the tube will be infinitely long so that the detected light will comprise parallel beams, *i.e.* it will be collimated light. In practice the tube, or its optical equivalent, need only be so long that a further increase in length would make very

little difference to the measured value of $L(\lambda)$. From eqns.4 and 5 we may write :

$$L(\lambda) = \frac{d^2 F(\lambda)}{dA \cdot d\Omega} = \frac{dI(\lambda)}{dA} \quad (\text{eqn.4})$$

Irradiance, $E(\lambda)$ [Wm⁻², qs⁻¹m⁻², lux]

Simply, the irradiance, $E(\lambda)$, is the radiant flux incident on a surface divided by the area of that element :

$$E(\lambda) = \frac{dF(\lambda)}{dA} \quad (\text{eqn.5})$$

This is illustrated in Fig.3(a):

Fig.3 Vector (cosine) irradiance collectors

There are a number of forms of irradiance commonly used in measuring the underwater light field, defined as follows :

Downwelling vector irradiance, $Ed(\lambda)$, the flux incident per unit area measured on a horizontal, upward facing light collector (cosine collector) as illustrated in Fig.3(b).

Irradiance may be described mathematically as the integration of radiance over a given solid angle, $d\omega$, so :

$$Ed(\lambda) = \int_{\theta=-\pi}^{\theta=\pi} \int_{\phi=0}^{\phi=2\pi} L(\lambda) \cdot \cos\theta \cdot d\omega \quad (\text{eqn.6})$$

The cosine term in eqn.5 arises because the light arrives at the (cosine) collector at an angle. Consider the beam of light of radiant flux, F , in Fig.3(c). Measured at right angles to the beam, the irradiance would be : $dE_1 = dF/dA_1$. However, the irradiance measured

by the collector is:

$$dE_2 = dF/dA_2, \text{ and since } dA_1 = dA_2 \cdot \cos\theta \text{ then } dE_2 = dE_1 \cdot \cos\theta$$

It is for this reason that the collector is called a *cosine* collector, and that the measured irradiance is termed *vector* irradiance.

Upwelling vector irradiance, $E_u(\lambda)$, the equivalent upward flux of irradiance as illustrated in *Fig.3(d)*:

$$E_u(\lambda) = \int_{\theta=-\pi}^{\theta=\pi} \int_{\phi=0}^{\phi=2\pi} L(\lambda) \cdot |\cos\theta| \cdot d\omega \quad (\text{eqn.7})$$

The total scalar irradiance, $E_o(\lambda)$, the integral of a radiance distribution at a point over all directions about a point and is equivalent to the *total* energy flux, dF/dt :

$$E_o(\lambda) = \int_{4\pi} L(\lambda) \cdot d\omega \quad (\text{eqn.8})$$

Total scalar irradiance is illustrated in *Fig.3(e)*. Of course, we may also have **upwelling scalar irradiance, $E_{ou}(\lambda)$** and **downwelling scalar irradiance, $E_{od}(\lambda)$**

$$E_{ou}(\lambda) = \int_{-2\pi} L(\lambda) \cdot d\omega ; \quad E_{od}(\lambda) = \int_{2\pi} L(\lambda) \cdot d\omega \quad (\text{eqn.9})$$

Summary

The flux of propagating light is described by a number of terms including:

- the radiant energy, $Q(\lambda)$ [J or q = quanta]
- radiant flux, $F(\lambda)$ [W, qs^{-1} , lumen]
- radiant intensity, $I(\lambda)$ [Wsr^{-1} , $qs^{-1}sr^{-1}$, lumen sr^{-1}]
- irradiance, $E(\lambda)$ [Wm^{-2} , $qs^{-1}m^{-2}$, lux]. Irradiance may be upwelling, downwelling, scalar or vector

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